

Model Dependence of Nucleon Resonance Parameters for

$P_{11}(1440)$, $D_{13}(1520)$ and $S_{11}(1535)$

L. Tiator and S. Kamalov, Uni Mainz

with participation of:

R. Arndt, R. Workman, C. Bennhold and I. Aznauryan

- Introduction:

Result from BRAG 2001 on $P_{33}(1232)$

- Definition of Photon Couplings or Helicity Amplitudes:

$A_{1/2}$, $A_{3/2}$, $S_{1/2}$

- Comparison with Partial Wave Analyses from:

GWU/SAID, MAID,

CC/Bennhold and UI M, DR/Aznauryan

- Preliminary Conclusions

Result of BRAG 2001 on Multipole Analysis in the Delta Region with Benchmark Dataset

summarized by R. Davidson, Proc. NSTAR2001 in Mainz, p. 204

	M1	E2	E2/M1 (%)
RPI	286	-7.2	-2.55
GWU	281	-7.2	-2.57
HA	281	-6.6	-2.35
MAID	275	-5.3	-1.93
KY	280	-6.2	-2.24
AZ	278	-6.3	-2.28
OM	288	-7.8	-2.77
AVG	281.3 ± 4.5	-6.6 ± 0.8	-2.38 ± 0.27

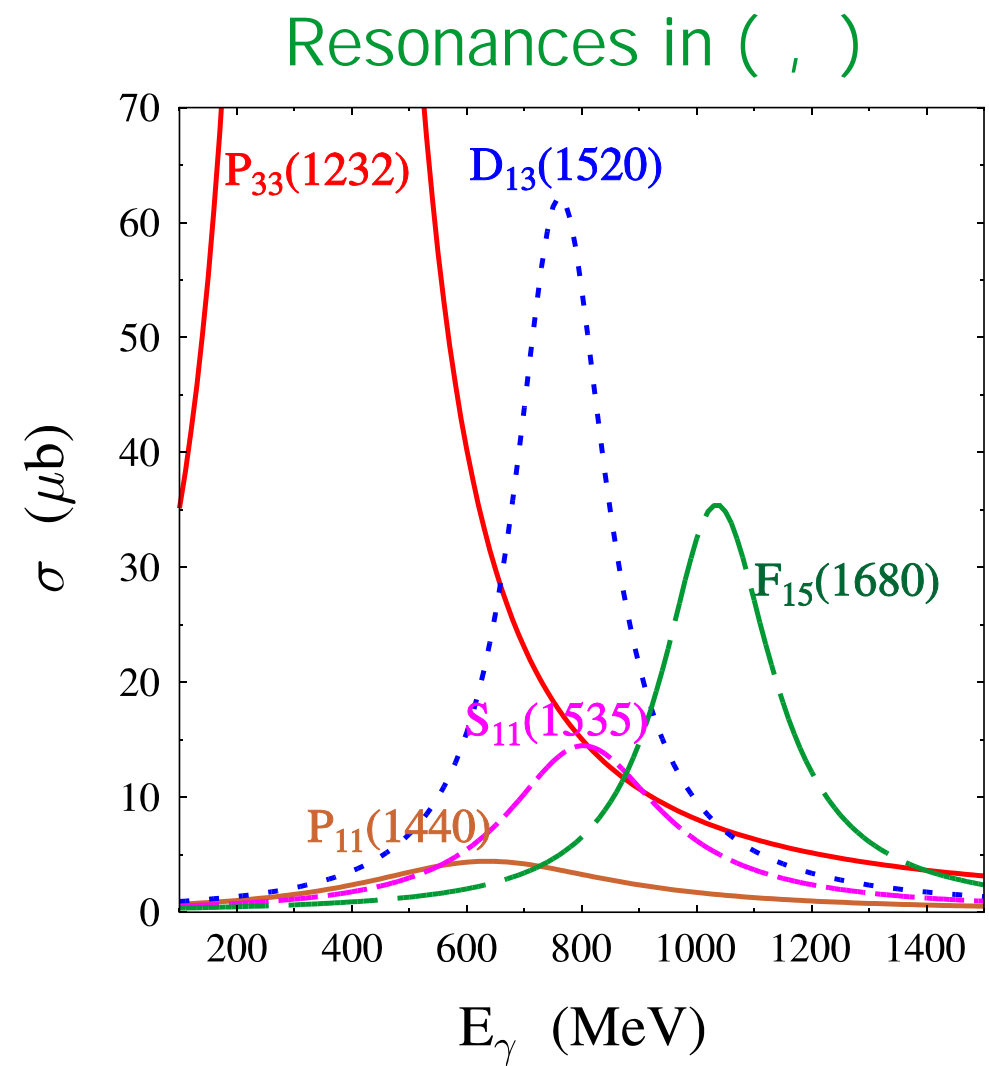


main conclusion:

model error in M1 coupling: $\sim 2\%$

model error in E2 and E/M $\sim 12\%$

Prominent Resonances in Pion Photoproduction



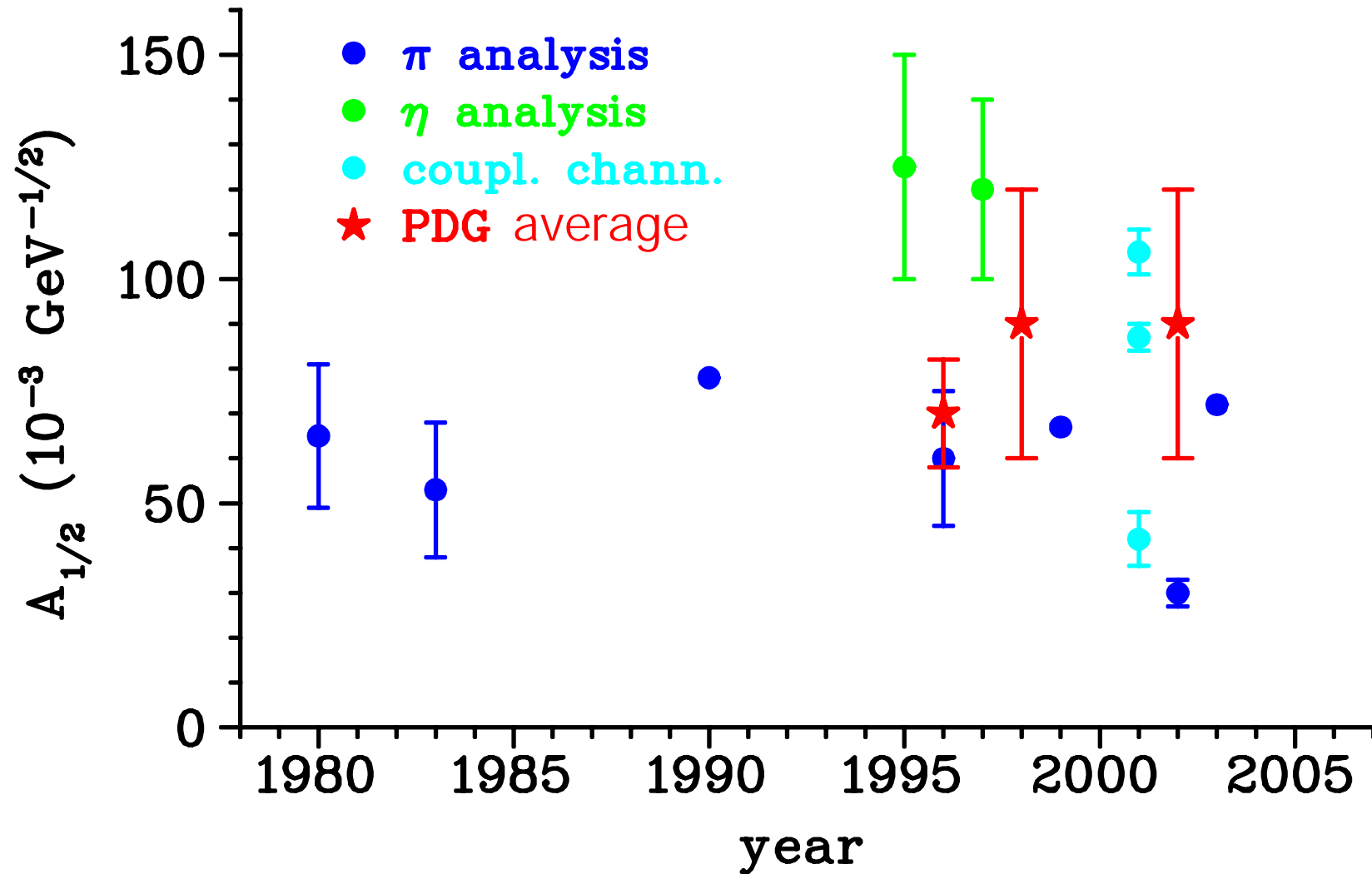
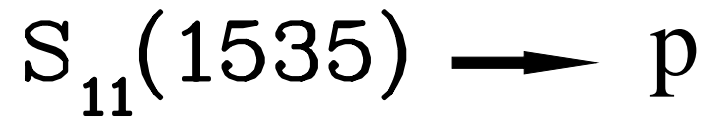
problems to face:

- we need very precise partial wave amplitudes
this is in principle possible, we are on the way
- we need to separate resonance and background
this is in principle impossible, but it could work approximately
- we need to know precisely the mass, width and single-pion branching ratio
this could be improved (in principle), very important for the width

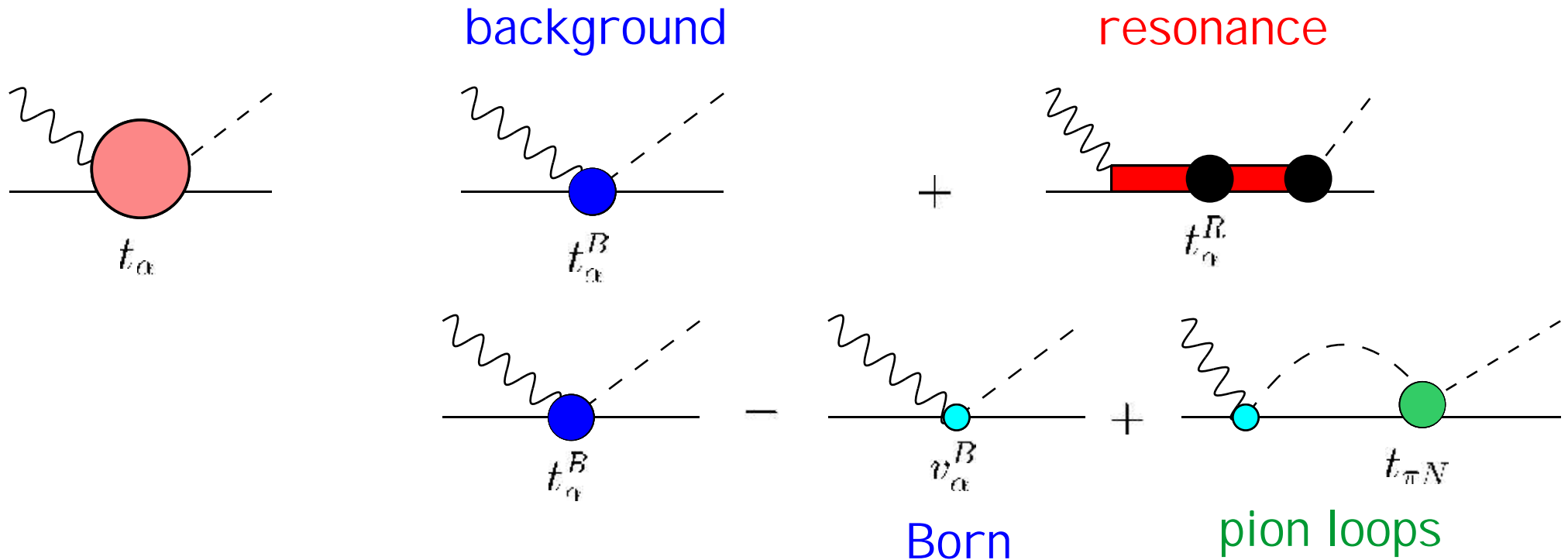
Particle Data Group 2002

	M_R	Γ_R	$\beta_\pi = \frac{\Gamma_\pi}{\Gamma_R}$	$pA_{1/2}$	$pA_{3/2}$
$P_{11}(1440)$	1440^{+30}_{-10}	350 ± 100	$.65 \pm .05$	-65 ± 4	—
$D_{13}(1520)$	1520^{+10}_{-5}	120^{+15}_{-10}	$.55 \pm .05$	-24 ± 9	166 ± 5
$S_{11}(1535)$	1535^{+20}_{-15}	150^{+100}_{-50}	$.45 \pm .10$	90 ± 30	—

Photon Decay Amplitudes



Dynamical Model Picture



the background is not only Born

the resonances are dressed

comparison with 3 methods for resonance - background separation

method a) used in SM02 by GWU/SAID group

b) used in previous SAID analyses

c) our (MAID) analysis

the partial wave amplitudes are fitted with a relatively simple form (~5-7 par.)
of smooth background plus Breit-Wigner resonance shape
in an energy region around the resonance position

$$a) \quad T_a = (1 + i t_{\pi N})(Born + A) + R t_{\pi N} + (C + i D)(\text{Im } t_{\pi N} - |t_{\pi N}|^2)$$

$$b) \quad T_b = [(1 + i t_{\pi N})(Born + A) + R t_{\pi N}] e^{i\phi}$$

$$c) \quad T_c = (1 + i t_{\pi N})(Born + A) + R t_{\pi N} e^{i\phi}$$

in details:

$$a) \quad T_a = (1 + i t_{\pi N})(Born + \mathbf{A}) + \mathbf{R} t_{\pi N} + (\mathbf{C} + i \mathbf{D})(\text{Im } t_{\pi N} - |t_{\pi N}|^2)$$

$$b) \quad T_b = [(1 + i t_{\pi N})(Born + \mathbf{A}) + \mathbf{R} t_{\pi N}] e^{i\phi}$$

$$c) \quad T_c = (1 + i t_{\pi N})(Born + \mathbf{A}) + \mathbf{R} t_{\pi N} e^{i\phi}$$

$$\text{Im } t_{\pi N} - |t_{\pi N}|^2 = (1 - \eta^2)/4 \quad (= 0, \text{ below } 2\pi \text{ threshold})$$

$$t_{\pi N} = \frac{M_R \Gamma_R}{M_R^2 - W^2 - i M_R \Gamma_{tot}(W)} \sqrt{\frac{\Gamma_\pi(W)}{\Gamma_\pi(M_R)}}$$

$\underset{W=M_R}{=} \underset{=}{i}$

energy dependent widths (including cusp effects)

$$\Gamma_{tot}(W) = \Gamma_\pi(W) + \Gamma_{2\pi}(W) + \Gamma_\eta(W)$$

$$\Gamma_{tot}(M_R) = \Gamma_R$$

up to 7 fit parameters:

$$\mathbf{A} = (a_1 + a_2 E)$$

$$\phi = (b_1 + b_2 E)(\text{Im } t_{\pi N} - |t_{\pi N}|^2)$$

$$\mathbf{R} = r_1$$

$$\mathbf{C} = c_1$$

$$\mathbf{D} = d_1$$

for P_{11} and S_{11} also mass and width:

$$M_R, \Gamma_R$$

reduced resonance multipoles:

$$\bar{\mathcal{A}}_\alpha(Q^2) = \frac{1}{c_{\pi N} f_{\pi R}(M_R)} \text{Im} \mathcal{A}_\alpha^{res}(M_R, Q^2)$$

kinematical factor

$$f_{\pi R}(M_R) = \left[\frac{1}{(2j+1)\pi} \frac{k_W m_N \Gamma_\pi}{|q| M_R \Gamma_R^2} \right]^{1/2}$$

isospin factor

$$c_{\pi N} = \begin{cases} -1/\sqrt{3} & : I = 1/2 \\ \sqrt{3}/2 & : I = 3/2 \end{cases}$$

photon couplings ($Q^2 = 0$):

$$\begin{aligned} S_{11} & : A_{1/2} = -\bar{E}_{0+} \\ P_{11} & : A_{1/2} = \bar{M}_{1-} \\ D_{13} & : A_{1/2} = -\frac{1}{2}(\bar{E}_{2-} - 3\bar{M}_{2-}) \\ & : A_{3/2} = -\frac{\sqrt{3}}{2}(\bar{E}_{2-} + \bar{M}_{2-}) \end{aligned}$$

e.g. for Roper resonance $P_{11}(1440)$ we get:

$$A_{1/2} = -\sqrt{\frac{6}{k^R} \frac{q^R M_R}{M_N}} \text{Im} M_{1-}^{res}(W = M_R)$$

Definition of Helicity Amplitudes or Photon Couplings

for resonance excitation: $\pi + N \rightarrow R$

$$A_{1/2} = -\sqrt{\frac{2}{k_w} \frac{f.s.}{k_w}} \langle R, \frac{1}{2} | J_+ | N, -\frac{1}{2} \rangle \quad A_{3/2} = -\sqrt{\frac{2}{k_w} \frac{f.s.}{k_w}} \langle R, \frac{3}{2} | J_+ | N, \frac{1}{2} \rangle$$

is a phase, which depends on the pion-decay matrix element
the photon couplings are **real** numbers

for $P_{11}(1440)$ in the simple harmonic oscillator CQM:

$$A_{1/2} = -\sqrt{\frac{\pi \alpha_{f.s.}}{3k_w} \frac{\mu_p}{6M_N} \alpha_0} \left(\frac{q^2}{\alpha_0^2}\right)^3 e^{-q^2/6\alpha_0^2} \quad \alpha_0 = 0.41 \text{ GeV} \text{ and } Q^2 = 0 \quad -24 \cdot 10^{-3} \text{ GeV}^{-1/2}$$

$$S_{11} : A_{1/2} = 168$$

$$D_{13} : A_{1/2} = -23$$

$$A_{3/2} = 125$$

Results of our Analysis

$$a) \quad T_a = (1 + i t_{\pi N})(Born + A) + R t_{\pi N} + (C + i D)(\text{Im } t_{\pi N} - |t_{\pi N}|^2)$$

$$b) \quad T_b = [(1 + i t_{\pi N})(Born + A) + R t_{\pi N}] e^{i\phi}$$

$$c) \quad T_c = (1 + i t_{\pi N})(Born + A) + R t_{\pi N} e^{i\phi}$$

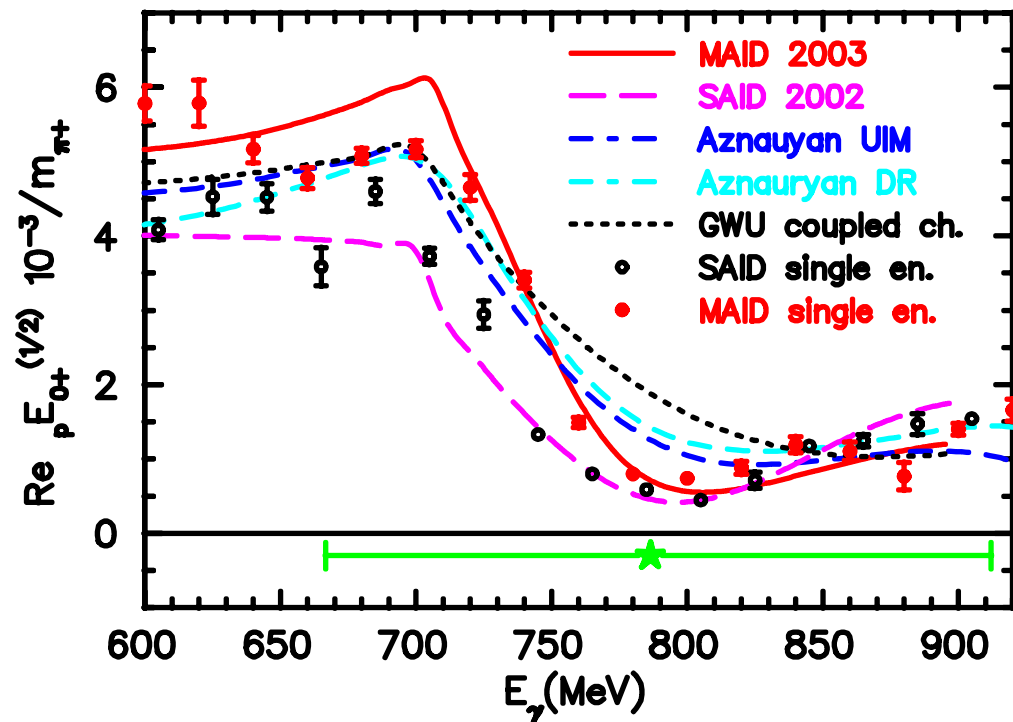
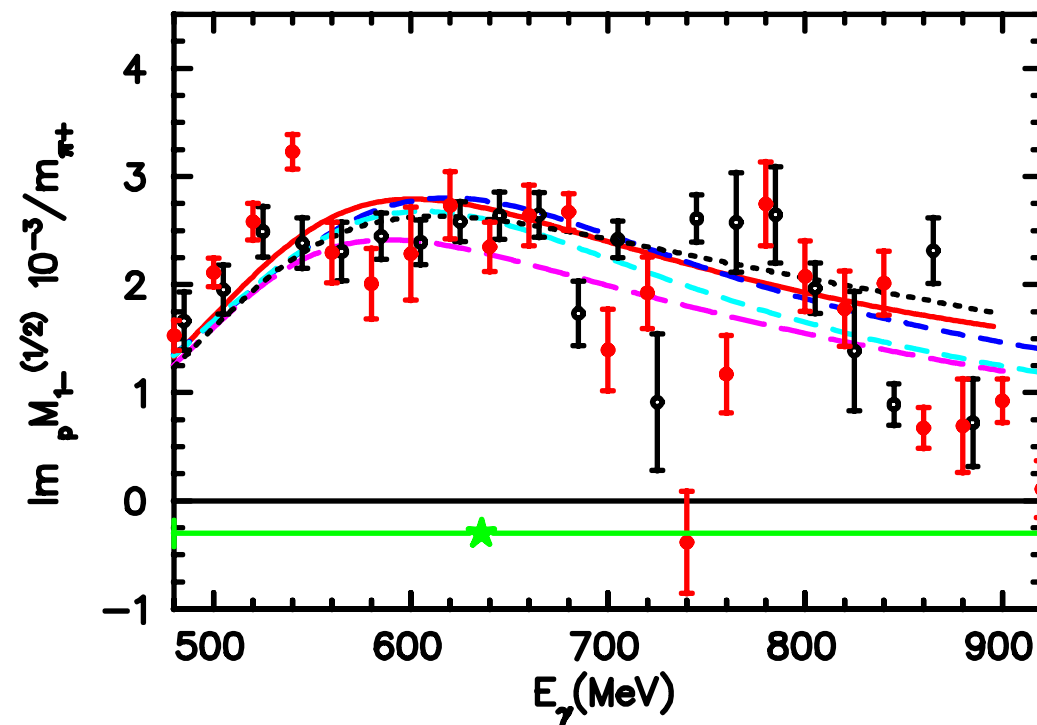
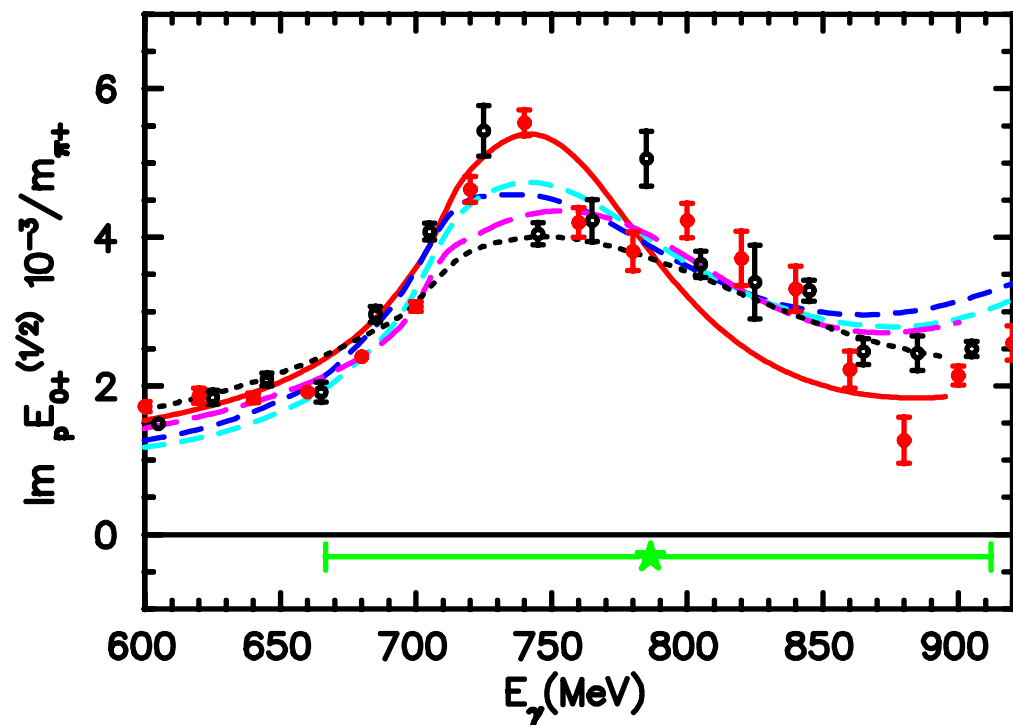
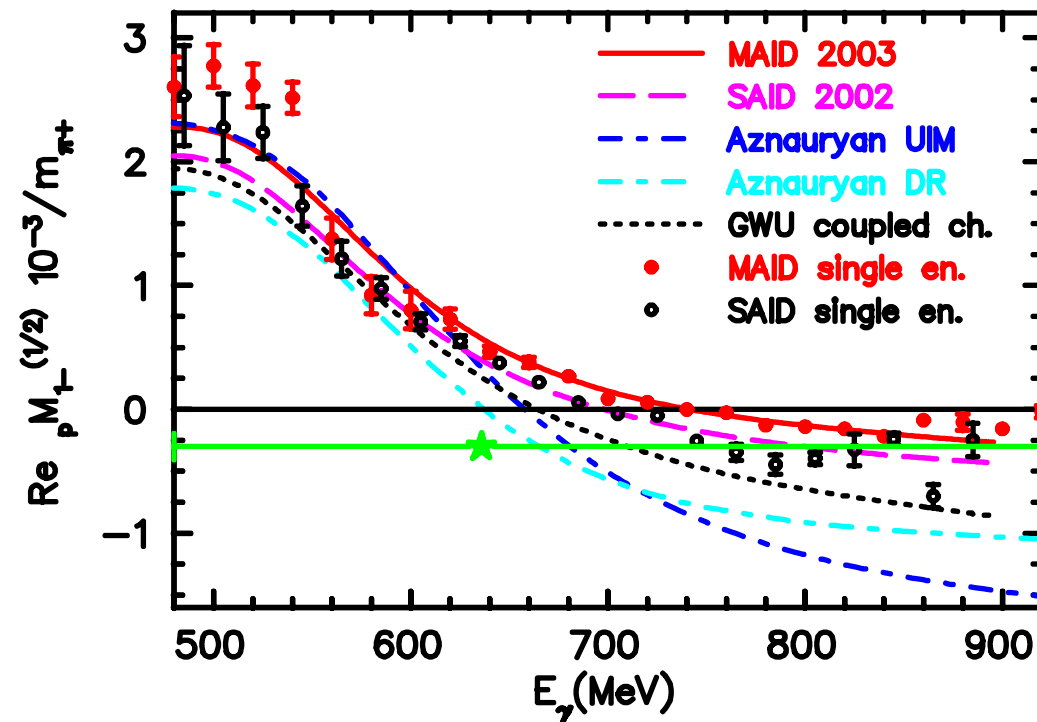
- most solutions can be fitted well with all 3 methods, only SAI D solutions for S11 is better fitted by method a)
- method b) and c) are very similar
we will show only results of method c) and compare it with a)
- most problematic case is the S11 partial wave
- less problematic, but with large experimental uncertainty is P11 pw
- no problems with the strong D13 pw

we have applied our 3 methods
to the partial wave analyses from different groups:

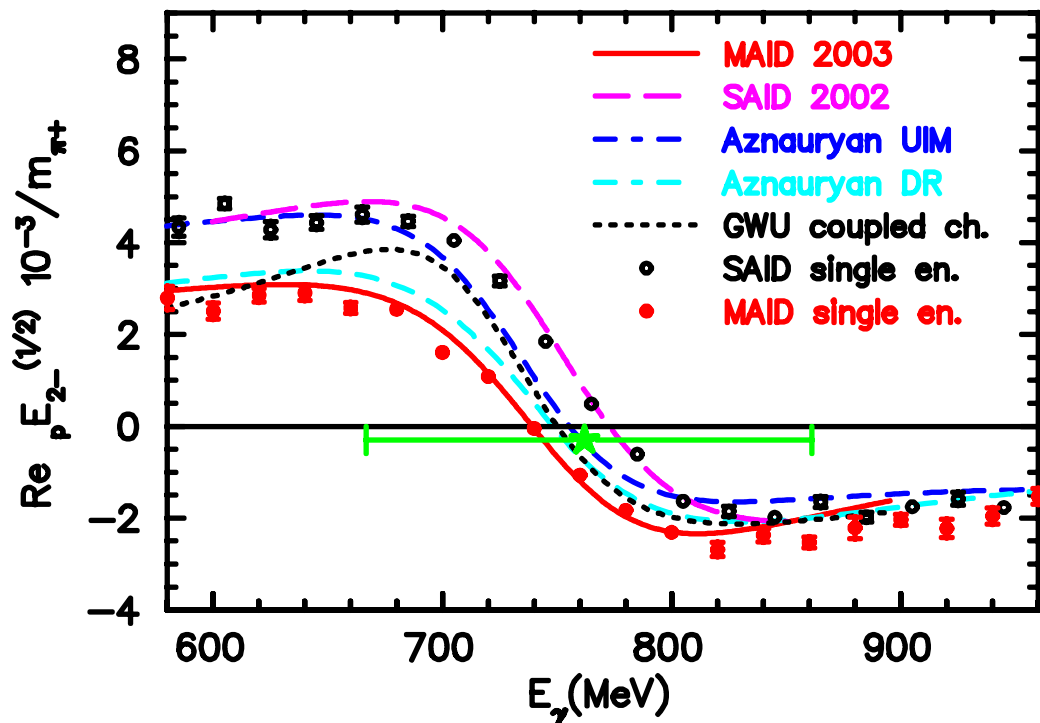
- MAI D global solution MAI D2003
- MAI D single-energy solution MAI Dse04
- SAI D/GWU global solution SM02
- SAI D/GWU single-energy solution SE02
- Bennhold/GWU coupled channels GWU/CC
- Aznauryan/JLab reggeized unitary isobar model AZ-UI M
- Aznauryan isobar model with dispersion relations AZ-DR

and we also compare to the following published results:

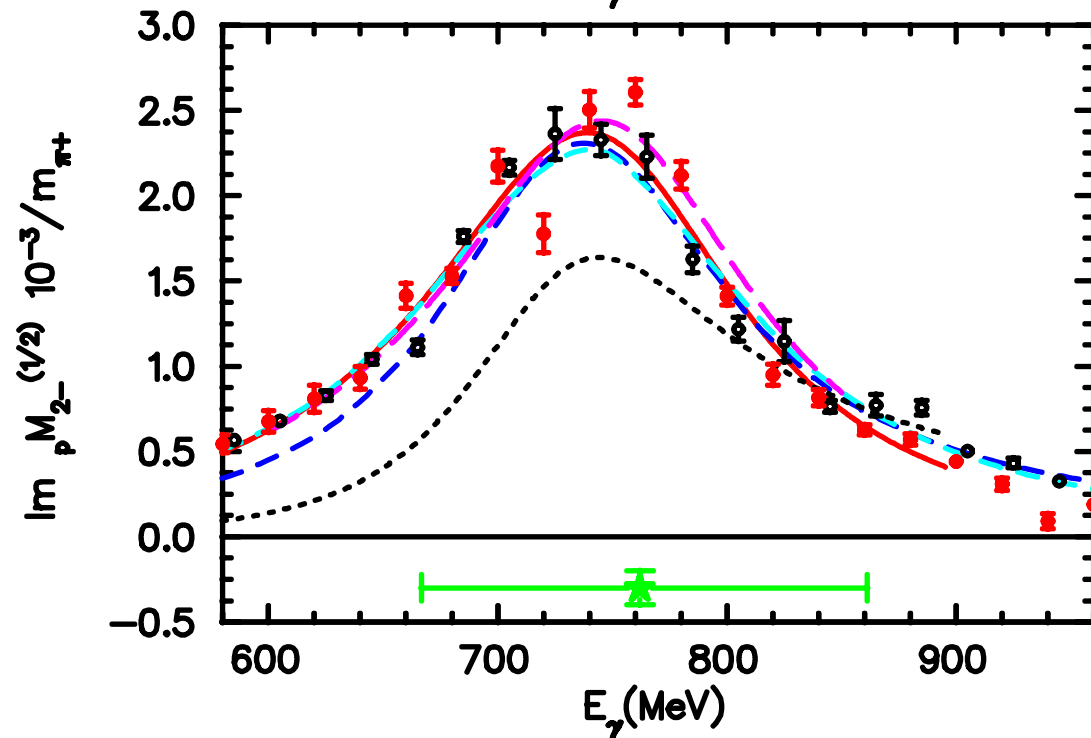
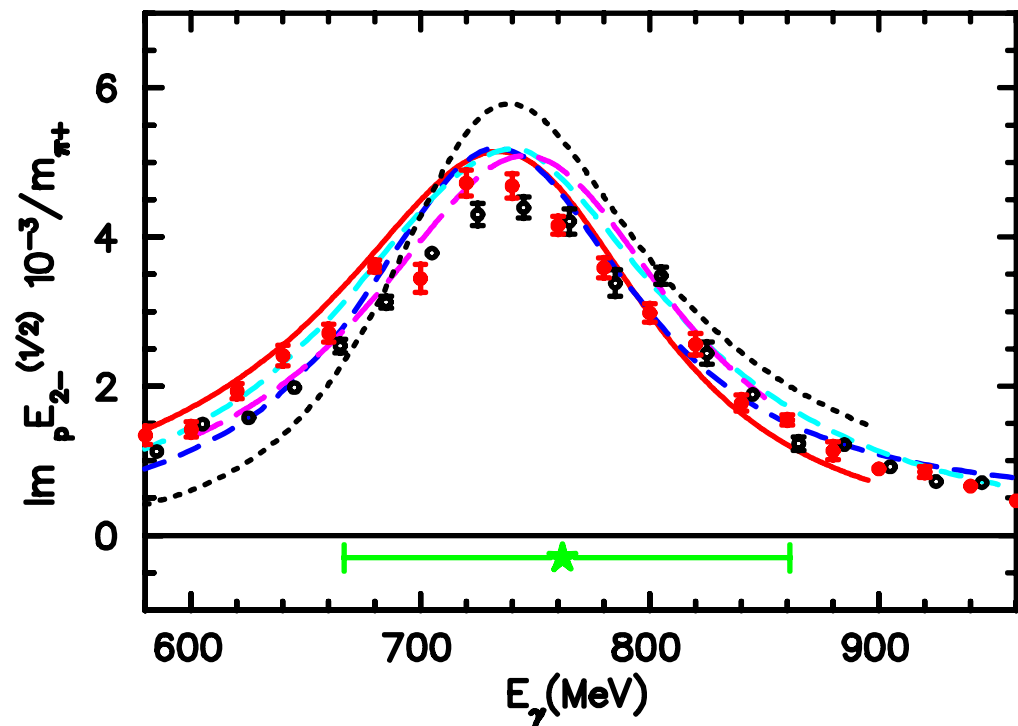
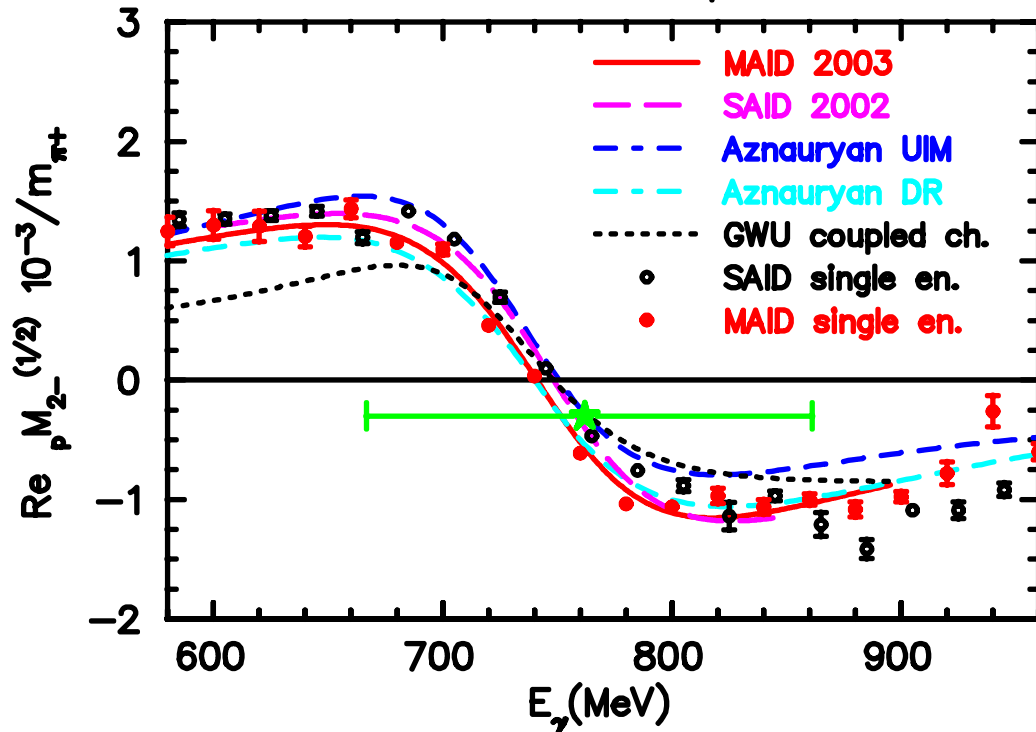
- PDG values
- GWU/SAI D analysis SM02, PR C 66,055213 (2002)
- , analysis for S11 by Krusche, Schadmand, review 2003

S_{11} partial wave ${}_p E_{0+}^{(1/2)}$  P_{11} partial wave ${}_p M_{1-}^{(1/2)}$ 

D_{13} partial wave ${}_p E_{2-}^{(1/2)}$



D_{13} partial wave ${}_p M_{2-}^{(1/2)}$



$S_{11}(1535)$

method *a*) $T_a = (1 + i t_{\pi N})(Born + A) + R t_{\pi N} + (C + i D)(Im t_{\pi N} - |t_{\pi N}|^2)$

method *c*) $T_c = (1 + i t_{\pi N})(Born + A) + R t_{\pi N} e^{i\phi}$

our analyses:



PDG



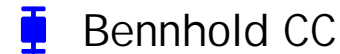
SM02



SAID



MAID



Bennhold CC



Az-UI M



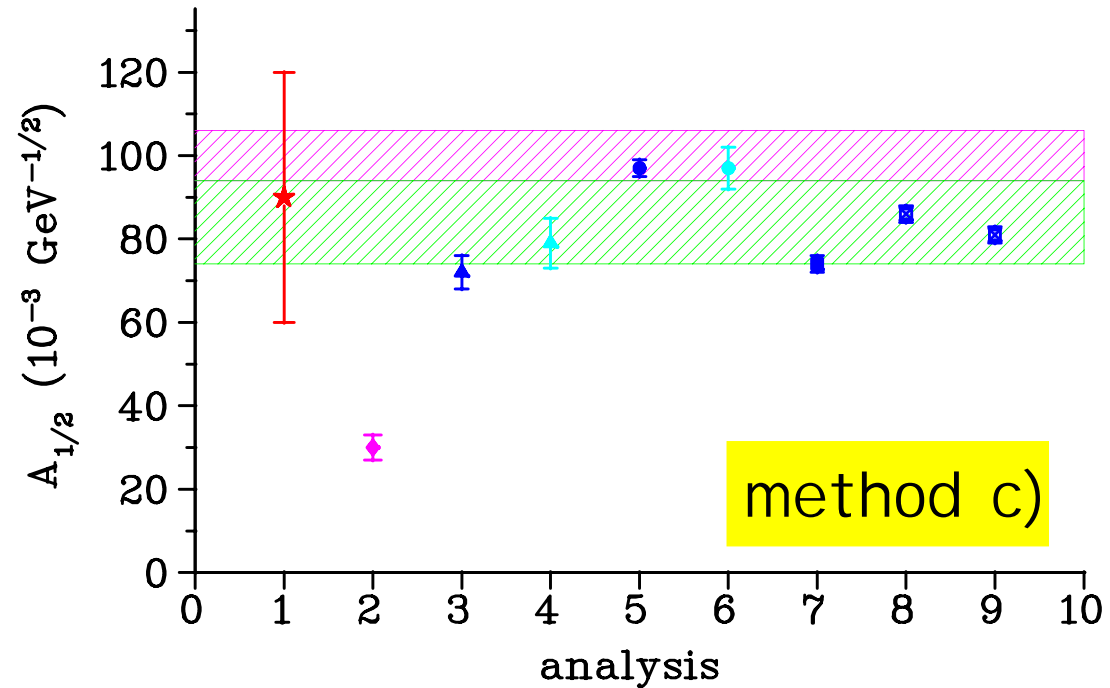
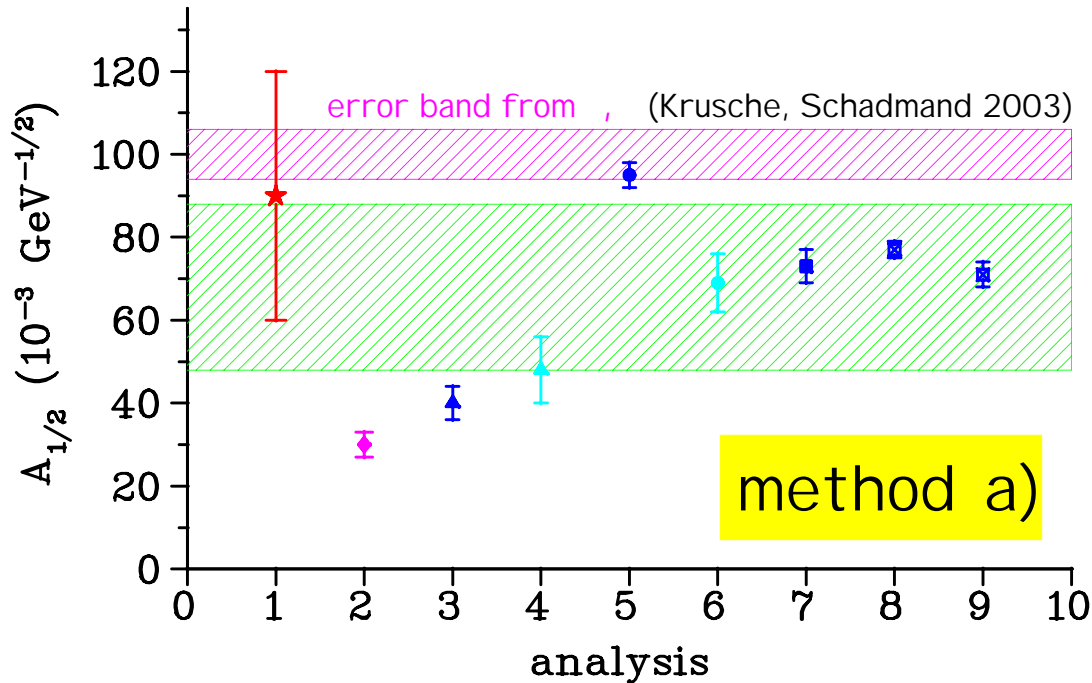
SAID s.e.



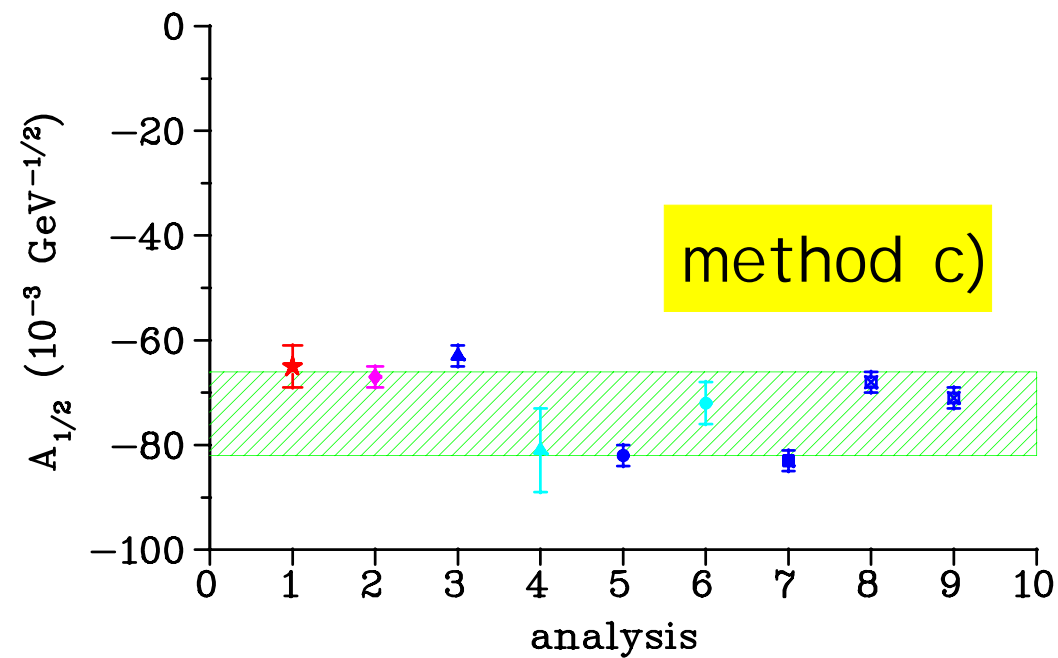
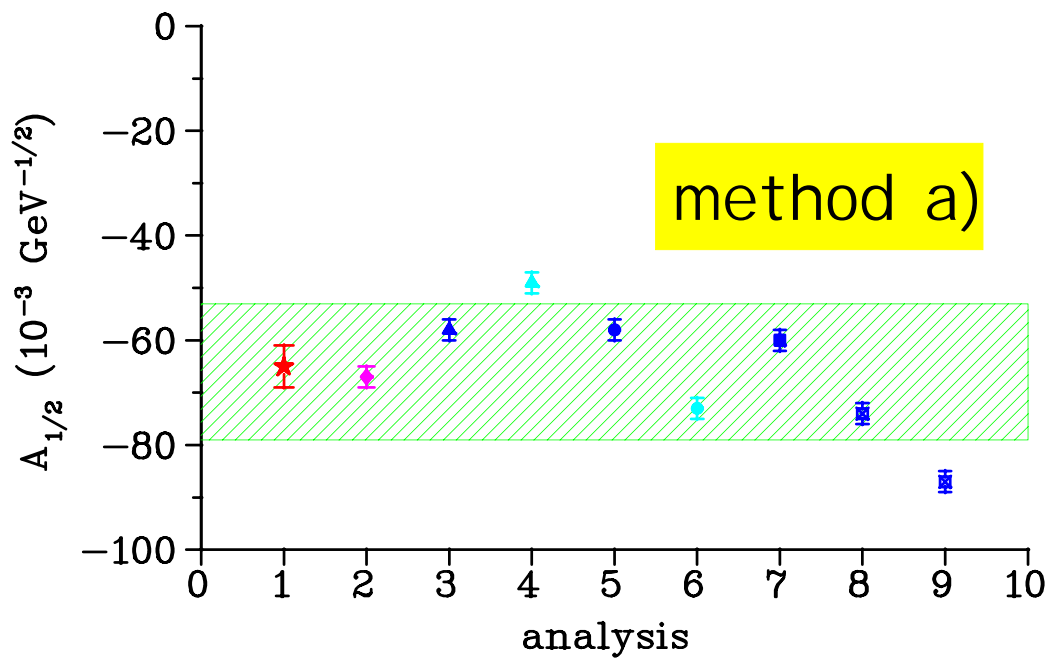
MAID s.e.



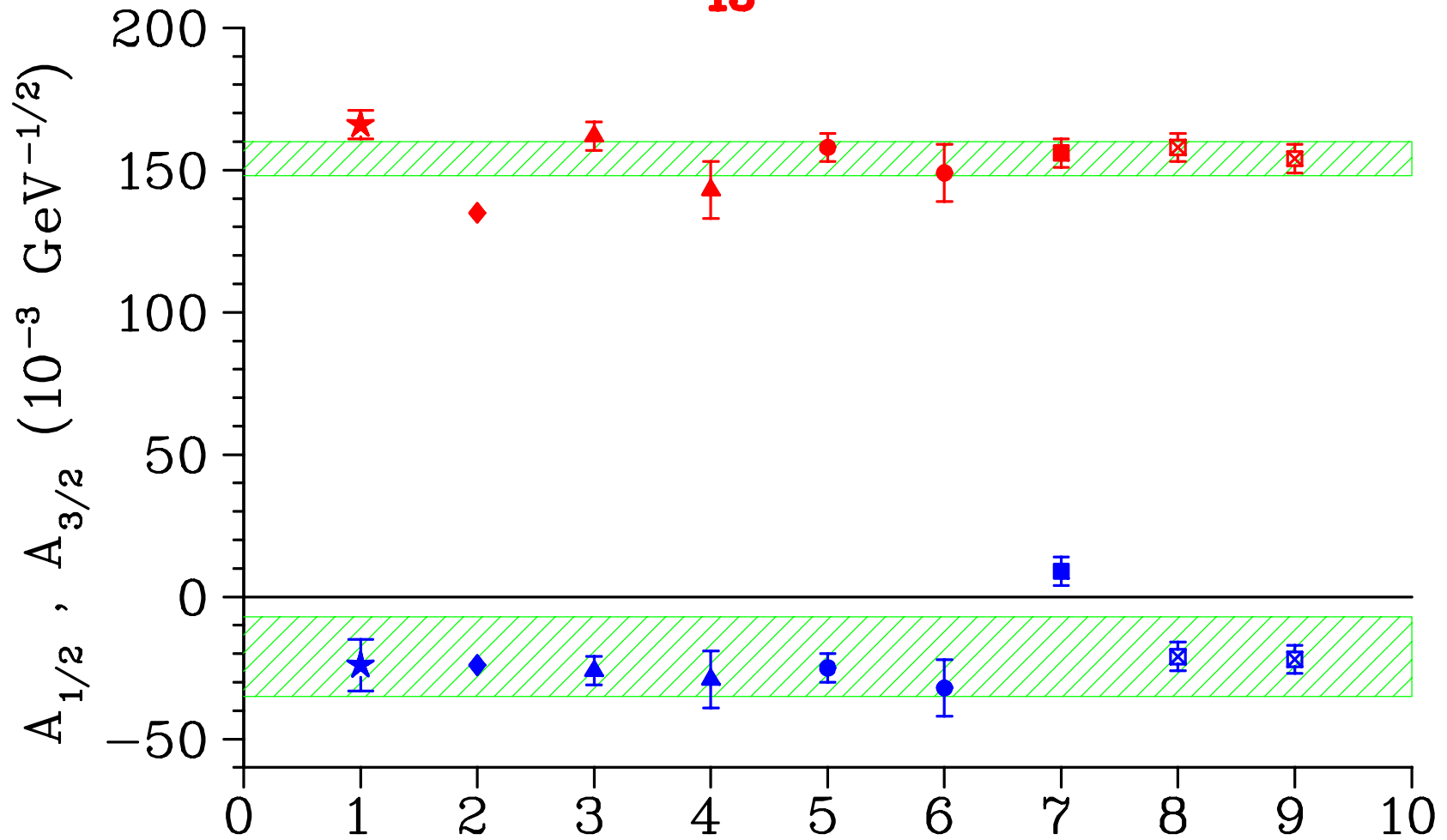
Az-DR



$P_{11}(1440)$



$D_{13}(1520)$



PDG

SAID

MAID

Bennhold CC

Az-DR

SM02

SAID s.e.

MAID s.e.

Az-UI M

(Preliminary) Conclusions

- method a) is problematic, even if it fits some SAID multipoles better,
it can mix background and resonance
method b) is similar to c)
we generally propose method c)

$$T_c = (1 + i t_{\pi N})(Born + A) + R t_{\pi N} e^{i\phi}$$

- it is a good idea to use normalized photon couplings with average values for M_R, Γ_R, β_π from PDG
- remaining uncertainties are:
 - a) from experimental data
 - b) from pw analysis
 - c) from energy range used in the BW fits